

Computer algebra application to numerical solving of nonlinear KdV-like equations

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Talk at ACA 2012, June 28, Sofia, Bulgaria

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Introduction

We apply our algorithmic approach [1][Gerdt, Blinkov, Mozzhilkin'2006] to generate finite-difference approximations (FDAs) to the nonlinear Korteweg-de Vries (KdV)-type equations and to construct numerical solutions.

The main steps of the approach are:

- 1 For a PDE system its completion to involution (Thomas decomposition, in general).
- 2 Conversion into the integral form of conservation law.
- 3 Enlargement with integral relations between unknown functions and their partial derivatives.
- 4 Discretization by applying the finite volume method and numerical integration.
- 5 Difference elimination of the partial derivatives to obtain FDA.
- 6 Verification of s(trong)-consistency [3][Gerdt'2012] of FDA.
- 7 Addition of discrete boundary condition (BVP) or/and initial conditions (IVP) and numerical solving.

Conservation Law Form for Two Independent Variables

A wide class of PDEs can be written in the **conservation law form**

$$\frac{\partial}{\partial x} \mathbf{F}(\mathbf{v}) + \frac{\partial}{\partial y} \mathbf{G}(\mathbf{v}) = 0.$$

Here \mathbf{v} is a m -vector function in unknown n -vector function \mathbf{u} and its partial derivatives $\mathbf{u}_x, \mathbf{u}_y, \mathbf{u}_{xx} \dots$; \mathbf{F} and \mathbf{G} are functions that map R^m into R^m .

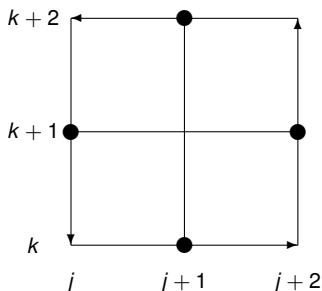
By **Green's theorem** (curl theorem in the plane), the above equation is equivalent to

$$\oint_{\Gamma} -\mathbf{G}(\mathbf{v}) dx + \mathbf{F}(\mathbf{v}) dy = 0.$$

where Γ is **arbitrary** closed contour.

Discretization

We set $\mathbf{u}(x, y) = \mathbf{u}(x_j, y_k) \equiv \mathbf{u}_{jk}$, $\mathbf{u}_x(x, y) = \mathbf{u}_x(x_j, y_k) \equiv (\mathbf{u}_x)_{jk}, \dots$, choose the integration contour, e.g.,



and add the relations for partial derivatives \mathbf{u}_x , \mathbf{u}_y , \mathbf{u}_{xx}, \dots

$$\int_{x_j}^{x_{j+2}} \mathbf{u}_x dx = \mathbf{u}(x_{j+2}, y) - \mathbf{u}(x_j, y), \quad \int_{y_k}^{y_{k+2}} \mathbf{u}_y dy = \mathbf{u}(x, y_{k+2}) - \mathbf{u}(x, y_k), \dots$$

Difference Elimination

Using a [numerical integration](#) method, e.g. the midpoint one, with

$$x_{j+1} - x_j = h_1, \quad y_{k+1} - y_k = h_2,$$

we rewrite the equations and the relations as

$$\begin{aligned} -(\mathbf{G}(\mathbf{v})_{j+1k} - \mathbf{G}(\mathbf{v})_{j+1k+2})h_1 + (\mathbf{F}(\mathbf{v})_{j+2k+1} - \mathbf{F}(\mathbf{v})_{jk+1})h_2 &= 0, \\ (\mathbf{u}_x)_{j+1k} \cdot 2h_1 &= \mathbf{u}_{j+2k} - \mathbf{u}_{jk}, \\ (\mathbf{u}_y)_{jk+1} \cdot 2h_2 &= \mathbf{u}_{jk+2} - \mathbf{u}_{jk}, \\ \dots\dots\dots \end{aligned}$$

A [conservative finite-difference approximation](#) to the PDEs is obtained by elimination of all partial derivatives $\mathbf{u}_x, \mathbf{u}_y, \mathbf{u}_{xx}, \dots$

KdV Equation

Describes waves on shallow water surfaces. It is example of an exactly solvable model.

$$u = u(t, x)$$

$$\frac{\partial}{\partial t}u + 6u\frac{\partial}{\partial x}u + \frac{\partial^3}{\partial x^3}u = 0.$$

Exact solution:

$$\theta = kx - \omega t, \quad k, \omega \text{ are constants,}$$

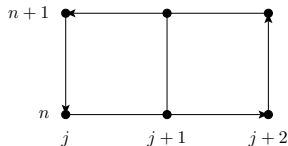
$$u = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(\theta).$$

Hereafter, we choose the orthogonal and uniform grid with the spacings τ in t and h in x .

Discretization

$$u_t + 6uu_x + u_{xxx} = 0,$$

$$\oint_{\partial\Omega} -(3u^2 + u_{xx}) dt + u dx = 0.$$



$$\int_{x_j}^{x_{j+1}} u_x dx = u(t, x_{j+1}) - u(t, x_j), \quad \int_{x_j}^{x_{j+1}} u_{xx} dx = u_x(t, x_{j+1}) - u_x(t, x_j),$$

$$\left\{ \begin{array}{l} -(1 + \theta_t - \theta_x^2 - \theta_t \theta_x^2) \circ (3u^2 + u_{xx}) \cdot \frac{\tau}{2} + (\theta_x \theta_t - \theta_x) \circ u \cdot 2h = 0, \\ \text{midpoint rule for } x \text{ and trapezoidal rule for } t \\ (\theta_x + 1) \circ u_x \cdot \frac{h}{2} = (\theta_x - 1) \circ u, \\ \text{trapezoidal rule} \\ \theta_x \circ u_{xx} \cdot 2h = (\theta_x^2 - 1) \circ u_x, \\ \text{midpoint rule} \end{array} \right.$$

θ_x and θ_t are the **right-shift operators** in x and t , respectively.

Input of LDA

To eliminate partial derivatives we replace $3u^2$ by F and use Maple package LDA (Linear Difference Algebra) [2][Gerdt, Robertz'2012]

```
> restart;
> libname:=libname, "/usr/local/lib/LDA":
> with(LDA):
> L:= [(-F(n,j)+F(n+1,j)-F(n,j+2)-F(n+1,j+2)) -
> (uxx(n,j)+uxx(n+1,j)-uxx(n,j+2)-uxx(n+1,j+2)))*tau/2+
> (u(n+1,j+1)-u(n,j+1))*2*h,
> (ux(n,j+1)+ux(n,j))*h/2-(u(n,j+1)-u(n,j)),
> 2*uxx(n,j+1)*h-(ux(n,j+2)-ux(n,j))];
```

$$\begin{aligned}
 L := & [(\frac{1}{2}(-F(n,j) - F(n+1,j) + F(n,j+2) + F(n+1,j+2) \\
 & - u_{xx}(n,j) - u_{xx}(n+1,j) + u_{xx}(n,j+2) + u_{xx}(n+1,j+2)))\tau \\
 & + (2(u(n+1,j+1) - u(n,j+1)))h, \\
 & (\frac{1}{2}(ux(n,j+1) + ux(n,j)))h - u(n,j+1) + u(n,j), \\
 & 2uxx(n,j+1)h - ux(n,j+2) + ux(n,j)]
 \end{aligned}$$

Output of LDA

> JanetBasis(L, [n,j], [uxx,ux,u,F],2);

$$\begin{aligned}
 & [[2\tau u(n, j + 1) - \tau u(n, j) + \tau u(n + 1, j + 4) - 2\tau u(n + 1, j + 3) + 4h^3 u(n + 1, j + 2) \\
 & + 2\tau u(n + 1, j + 1) - \tau u(n + 1, j) + \tau u(n, j + 4) - 2\tau u(n, j + 3) \\
 & - 4h^3 u(n, j + 2) + h^2 \tau F(n + 1, j + 3) - h^2 \tau F(n + 1, j + 1) + h^2 \tau F(n, j + 3) \\
 & - h^2 \tau F(n, j + 1), hux(n, j + 1) + hux(n, j) - 2u(n, j + 1) + 2u(n, j), \\
 & h^2 uxx(n, j + 1) - u(n, j + 2) + 2u(n, j + 1) - u(n, j), -h^2 \tau uxx(n + 1, j) - \\
 & h^2 \tau uxx(n, j) + \tau u(n + 1, j + 3) - 2\tau u(n + 1, j + 2) + (\tau + 4h^3)u(n + 1, j + 1) \\
 & + \tau u(n, j + 3) - 2\tau u(n, j + 2) + (\tau - 4h^3)u(n, j + 1) + h^2 \tau F(n + 1, j + 2) \\
 & - h^2 \tau F(n + 1, j) + h^2 \tau F(n, j + 2) - h^2 \tau F(n, j)], [n, j], [uxx, ux, u, F]]
 \end{aligned}$$

Extraction of FDA

> collect(%[1,1]/(4*tau*h**3),[tau,h,sigma]);

$$\frac{\frac{1}{4}F(n+1,j+3) - \frac{1}{4}F(n+1,j+1) + \frac{1}{4}F(n,j+3) - \frac{1}{4}F(n,j+1)}{h}$$

$$+ (\frac{1}{4}u(n+1,j+4) - \frac{1}{2}u(n+1,j+3) + \frac{1}{2}u(n,j+1) - \frac{1}{4}u(n,j)$$

$$- \frac{1}{4}u(n+1,j) + \frac{1}{4}u(n,j+4) - \frac{1}{2}u(n,j+3) + \frac{1}{2}u(n+1,j+1))/h^3$$

$$+ \frac{u(n+1,j+2) - u(n,j+2)}{\tau}$$

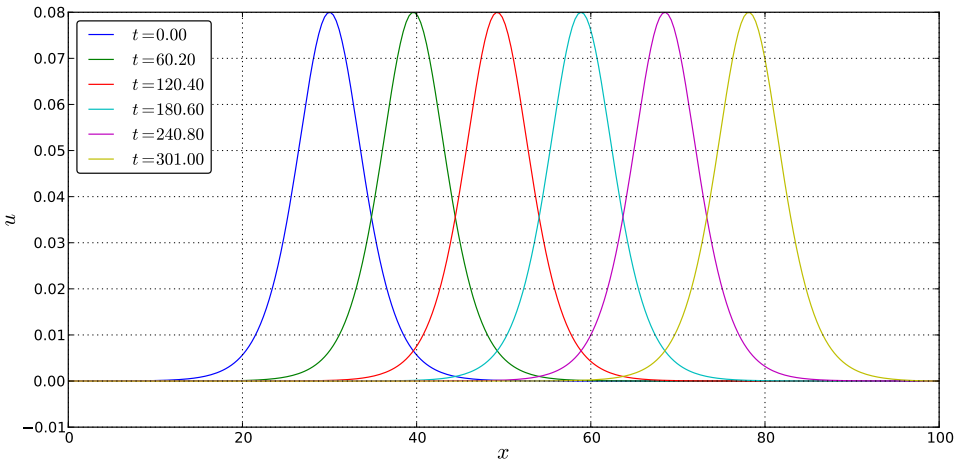
Finite Difference Approximation

Substituting back $F = 3u^2$ we obtain

$$\begin{aligned} & \frac{u_j^{n+1} - u_j^n}{\tau} + \\ & + 3 \frac{(u_{j+1}^{2n+1} - u_{j-1}^{2n+1}) + (u_{j+1}^{2n} - u_{j-1}^{2n})}{4h} + \\ & + \frac{(u_{j+2}^{n+1} - 2u_{j+1}^{n+1} + 2u_{j-1}^{n+1} - u_{j-2}^{n+1}) + (u_{j+2}^n - 2u_{j+1}^n + 2u_{j-1}^n - u_{j-2}^n)}{4h^3} = 0. \end{aligned}$$

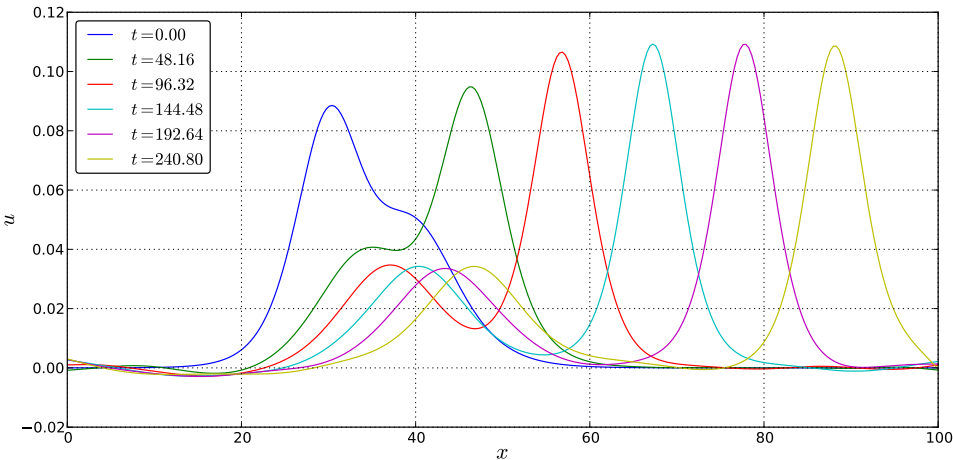
The l.h.s. is a [difference Gröbner basis](#) (no nontrivial \mathcal{S} -polynomials), and by the construction it is consistent with KdV in the limit $h \rightarrow 0, \tau \rightarrow 0$.

Therefore [FDA is s-consistent](#) [3][Gerdt'2012].



$$\text{IVP (soliton solution)} \quad f(t, x, k, \omega) = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(kx - \omega t)$$

$$u(0, x) = f(0, x + 30, 0.2, 4 \cdot 0.2^3)$$



Double soliton solution with $f(t, x, k, \omega) = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(kx - \omega t)$

$$u(0, x) = f(0, x + 30, 0.2, 4 \cdot 0.2^3) + f(0, x + 40, 0.15, 4 \cdot 0.15^3)$$

KdV-Like Equations

Now we consider the following 3-parametric family of PDEs

$$\frac{\partial}{\partial t} u + 6u \frac{\partial}{\partial x} u + \frac{\partial^3}{\partial x^3} u - s_1 u^2 \frac{\partial}{\partial x} u - s_2 \frac{\partial^2}{\partial x^2} u - su = 0,$$

s, s_1, s_2 are constants.

This family describes propagation of nonlinear deformation waves in elastic cylinder shells containing viscous incompressible liquid [?][Mogilevich, Blinkova, Ivanov'2011].

$s = s_1 = s_2 = 0 \implies$ KdV equation.

The sign of s characterizes the shell material:

- $s > 0$ – nonorganic,
- $s < 0$ – living organisms,
- $s = 0$ – rubber.

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One-Parametric Family

$$s = s_1 = 0, s_2 \neq 0$$

$$\text{equation: } \frac{\partial}{\partial t} u + 6u \frac{\partial}{\partial x} u + \frac{\partial^3}{\partial x^3} u - s_2 \frac{\partial^2}{\partial x^2} u = 0,$$

exact solutions:

$$\theta = \frac{1}{10} s_2 x - \omega t,$$

$$u = \frac{5}{3} \frac{\omega}{s_2} + \frac{1}{50} s_2^2 - \frac{1}{25} s_2^2 \tanh(\theta) - \frac{1}{50} s_2^2 \tanh^2(\theta),$$

and

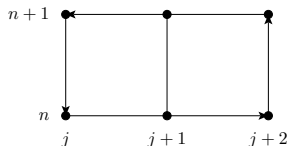
$$\theta = -\frac{1}{10} s_2 x - \omega t,$$

$$u = -\frac{5}{3} \frac{\omega}{s_2} + \frac{1}{50} s_2^2 + \frac{1}{25} s_2^2 \tanh(\theta) - \frac{1}{50} s_2^2 \tanh^2(\theta).$$

Discretization

$$u_t + 6uu_x + u_{xxx} - s_2 u_{xx} = 0,$$

$$\oint_{\partial\Omega} -(3u^2 + u_{xx} - s_2 u_x) dt + u dx = 0,$$



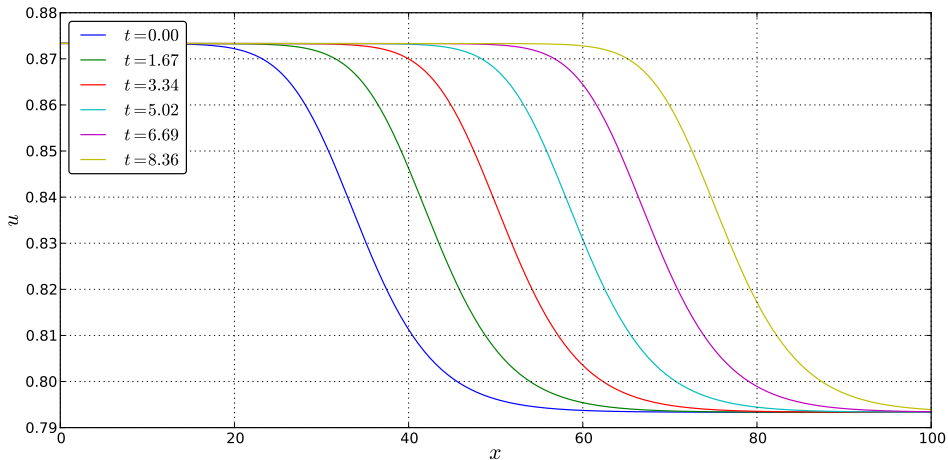
$$\int_{x_j}^{x_{j+1}} u_x dx = u(t, x_{j+1}) - u(t, x_j), \quad \int_{x_j}^{x_{j+1}} u_{xx} dx = u_x(t, x_{j+1}) - u_x(t, x_j).$$

$$\left\{ \begin{array}{l} -(1 + \theta_t - \theta_x^2 - \theta_t \theta_x^2) \circ (3u^2 + u_{xx} - s_2 u_x) \cdot \frac{\tau}{2} + \\ \quad + (\theta_x \theta_t - \theta_x) \circ u \cdot 2h = 0, \\ (\theta_x + 1) \circ u_x \cdot \frac{h}{2} = (\theta_x - 1) \circ u \\ \theta_x \circ u_{xx} \cdot 2h = (\theta_x^2 - 1) \circ u_x. \end{array} \right.$$

Finite Difference Approximation

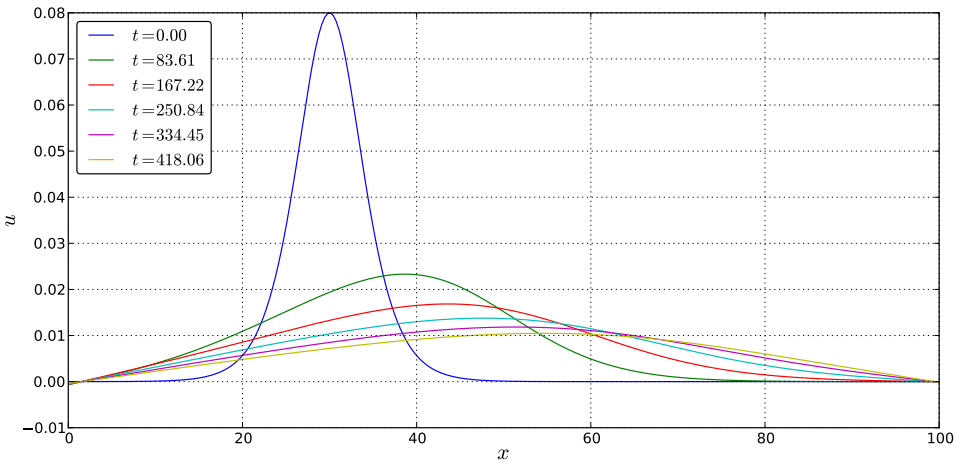
S-consistent FDA:

$$\begin{aligned}
 & \frac{u_j^{n+1} - u_j^n}{\tau} + \\
 & + 3 \frac{(u_{j+1}^{2n+1} - u_{j-1}^{2n+1}) + (u_{j+1}^{2n} - u_{j-1}^{2n})}{4h} + \\
 & + \frac{(u_{j+2}^{n+1} - 2u_{j+1}^{n+1} + 2u_{j-1}^{n+1} - u_{j-2}^{n+1}) + (u_{j+2}^n - 2u_{j+1}^n + 2u_{j-1}^n - u_{j-2}^n)}{4h^3} - \\
 & - s_2 \frac{(u_{j+1}^{n+1} - 2u_j^{n+1} + u_{j-1}^{n+1}) + (u_{j+1}^n - 2u_j^n + u_{j-1}^n)}{2h^2} = 0.
 \end{aligned}$$



Initial data $f(t, x, s_2, \omega) = \frac{5}{3} \frac{\omega}{s_2} + \frac{1}{50} s_2^2 -$
 $-\frac{1}{25} s_2^2 \tanh\left(\frac{1}{10} s_2 x - \omega t\right) - \frac{1}{50} s_2^2 \tanh^2\left(\frac{1}{10} s_2 x - \omega t\right)$

$$u(0, x) = f(0, x + 30, 1.0, 0.5)$$



Soliton dynamics (blooming) for IVP with $s_2 = 1.0$ and

$$f(t, x, k, \omega) = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(kx - \omega t)$$

$$u(0, x) = f(0, x + 30, 0.2, 4 \cdot 0.2^3)$$

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Two-Parametric Family

$$\text{equation: } \frac{\partial}{\partial t} u + 6u \frac{\partial}{\partial x} u + \frac{\partial^3}{\partial x^3} u - s_1 u^2 \frac{\partial}{\partial x} u - s_2 \frac{\partial^2}{\partial x^2} u = 0,$$

exact solution:

$$\theta = kx + t \left(-9 \frac{k}{s_1} + \frac{1}{6} k s_2^2 + 2k^3 \right),$$

$$u = \frac{3}{s_1} \pm \frac{s_2 \sqrt{6}}{6 \sqrt{s_1}} \pm \frac{k \sqrt{6}}{\sqrt{s_1}} \tanh(\theta).$$

one-parametric subfamily:

$$\frac{\partial}{\partial t} u + 6u \frac{\partial}{\partial x} u + \frac{\partial^3}{\partial x^3} u - \frac{54}{s_2^2} u^2 \frac{\partial}{\partial x} u - s_2 \frac{\partial^2}{\partial x^2} u = 0,$$

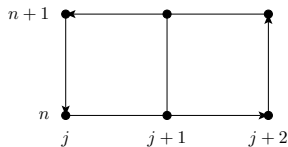
$$\theta = kx + 2tk^3,$$

$$u = -\frac{1}{3} k s_2 \tanh(\theta).$$

Discretization

$$u_t + 6uu_x + u_{xxx} - s_1 u^2 u_{xx} - s_2 u_{xx} = 0,$$

$$\oint_{\partial\Omega} -(3u^2 + u_{xx} - \frac{s_1}{3}u^3 - s_2 u_x) dt + u dx = 0,$$



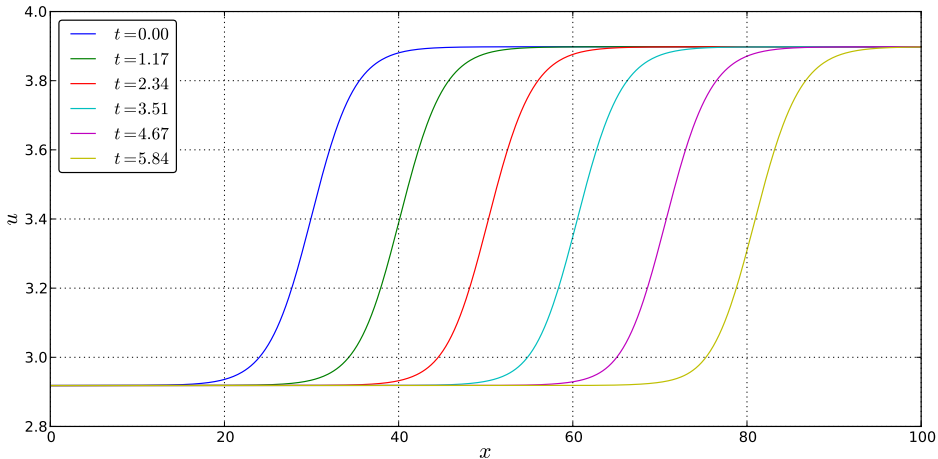
$$\int_{x_j}^{x_{j+1}} u_x dx = u(t, x_{j+1}) - u(t, x_j), \quad \int_{x_j}^{x_{j+1}} u_{xx} dx = u_x(t, x_{j+1}) - u_x(t, x_j).$$

$$\left\{ \begin{array}{l} -(1 + \theta_t - \theta_x^2 - \theta_t \theta_x^2) \circ (3u^2 + u_{xx} - \frac{s_1}{3}u^3 - s_2 u_x) \cdot \frac{\tau}{2} + \\ \quad + (\theta_x \theta_t - \theta_x) \circ u \cdot 2h = 0, \\ (\theta_x + 1) \circ u_x \cdot \frac{h}{2} = (\theta_x - 1) \circ u, \\ \theta_x \circ u_{xx} \cdot 2h = (\theta_x^2 - 1) \circ u_x. \end{array} \right.$$

Finite Difference Approximation

S-consistent FDA:

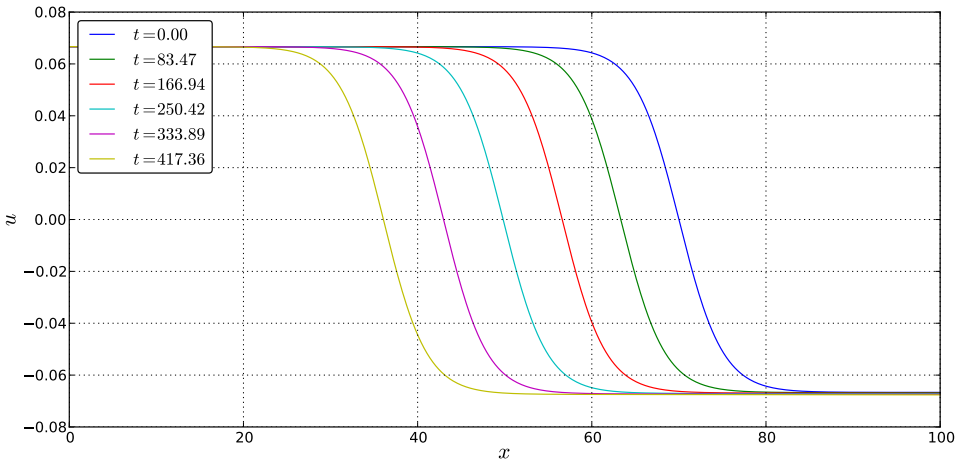
$$\begin{aligned}
 & \frac{u_j^{n+1} - u_j^n}{\tau} + \\
 & + 3 \frac{(u_{j+1}^{2n+1} - u_{j-1}^{2n+1}) + (u_{j+1}^{2n} - u_{j-1}^{2n})}{4h} + \\
 & + \frac{(u_{j+2}^{n+1} - 2u_{j+1}^{n+1} + 2u_{j-1}^{n+1} - u_{j-2}^{n+1}) + (u_{j+2}^n - 2u_{j+1}^n + 2u_{j-1}^n - u_{j-2}^n)}{4h^3} - \\
 & - \frac{s_1}{3} \frac{(u_{j+1}^{3n+1} - u_{j-1}^{3n+1}) + (u_{j+1}^{3n} - u_{j-1}^{3n})}{4h} - \\
 & - s_2 \frac{(u_{j+1}^{n+1} - 2u_j^{n+1} + u_{j-1}^{n+1}) + (u_{j+1}^n - 2u_j^n + u_{j-1}^n)}{2h^2} = 0.
 \end{aligned}$$



IVP for the exact solution with

$$f(t, x, k, s_1, s_2) = \frac{3}{s_1} + \frac{s_2 \sqrt{6}}{6\sqrt{s_1}} + \frac{k\sqrt{6}}{\sqrt{s_1}} \tanh \left(kx + t \left(-9\frac{k}{s_1} + \frac{1}{6}ks_2^2 + 2k^3 \right) \right)$$

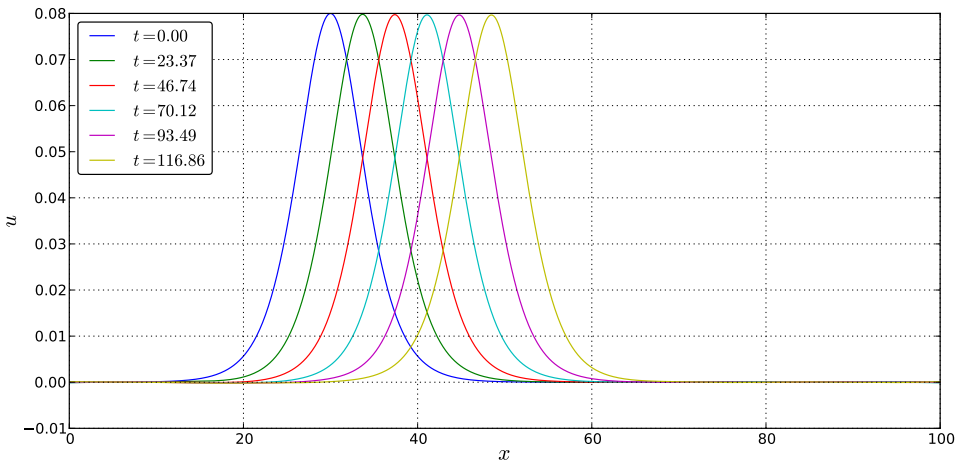
$$u(0, x) = f(0, x + 30, 0.2, 1.0, 1.0)$$



IVP for the exact solution with $s_1 = 54/s_2^2$ and

$$f(t, x, k, s_2) = -\frac{1}{3}ks_2 \tanh(kx + 2tk^3)$$

$$u(0, x) = f(0, x + 70, 0.2, 1.0)$$



Soliton dynamics for $s_1 = 1.0$, $s_2 = 0.0$ and

$$f(t, x, k, \omega) = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(kx - \omega t)$$

$$u(0, x) = f(0, x + 30, 0.2, 4 \cdot 0.2^3)$$

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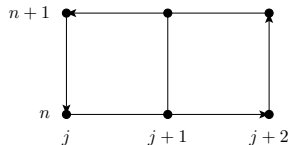
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Discretization

$$u_t + 6uu_x + u_{xxx} - s_1 u^2 u_{xx} - s_2 u_{xx} - su = 0,$$

$$\oint_{\partial\Omega} \left(3u^2 + u_{xx} - \frac{s_1}{3} u^3 - s_2 u_x \right) dt + u dx -$$

$$- s \iint_{\Omega} u dt dx = 0,$$



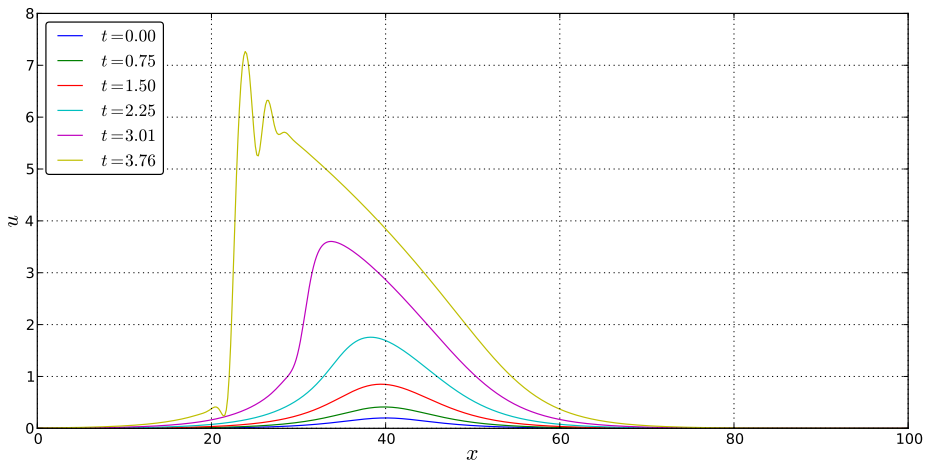
$$\int_{x_j}^{x_{j+1}} u_x dx = u(t, x_{j+1}) - u(t, x_j), \quad \int_{x_j}^{x_{j+1}} u_{xx} dx = u_x(t, x_{j+1}) - u_x(t, x_j).$$

$$\left\{ \begin{array}{l} -(1 + \theta_t - \theta_x^2 - \theta_t \theta_x^2) \circ (3u^2 + u_{xx} - \frac{s_1}{3} u^3 - s_2 u_x) \cdot \frac{\tau}{2} + \\ \quad + (\theta_x \theta_t - \theta_x) \circ u \cdot 2h - s(\theta_x \theta_t + \theta_x) \circ u \cdot h\tau = 0, \\ (\theta_x + 1) \circ u_x \cdot \frac{h}{2} = (\theta_x - 1) \circ u, \\ \theta_x \circ u_{xx} \cdot 2h = (\theta_x^2 - 1) \circ u_x. \end{array} \right.$$

Finite-Difference Approximation

S-consistent FDA:

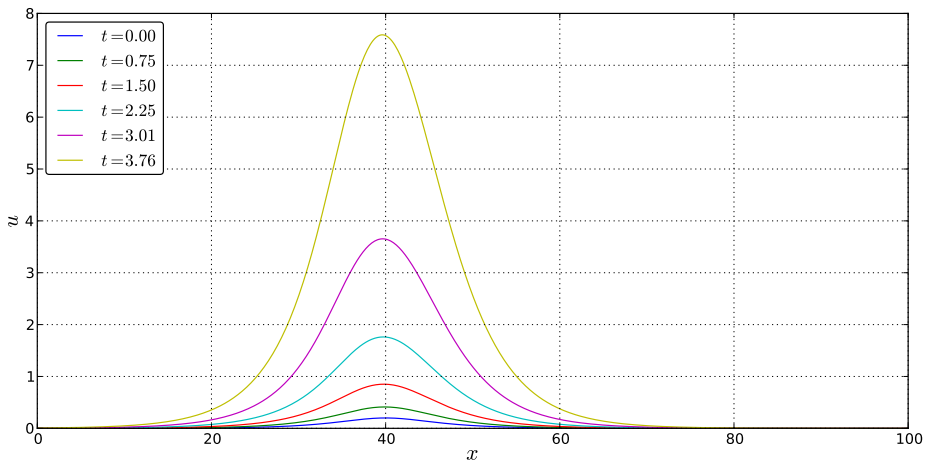
$$\begin{aligned}
 & \frac{u_j^{n+1} - u_j^n}{\tau} + \\
 & + 3 \frac{(u_{j+1}^{2n+1} - u_{j-1}^{2n+1}) + (u_{j+1}^{2n} - u_{j-1}^{2n})}{4h} + \\
 & + \frac{(u_{j+2}^{n+1} - 2u_{j+1}^{n+1} + 2u_{j-1}^{n+1} - u_{j-2}^{n+1}) + (u_{j+2}^n - 2u_{j+1}^n + 2u_{j-1}^n - u_{j-2}^n)}{4h^3} - \\
 & - \frac{s_1 (u_{j+1}^{3n+1} - u_{j-1}^{3n+1}) + (u_{j+1}^{3n} - u_{j-1}^{3n})}{3 \cdot 4h} - \\
 & - s_2 \frac{(u_{j+1}^{n+1} - 2u_j^{n+1} + u_{j-1}^{n+1}) + (u_{j+1}^n - 2u_j^n + u_{j-1}^n)}{2h^2} - \\
 & - s(u_j^{n+1} + u_j^n) \frac{1}{2} = 0.
 \end{aligned}$$



Soliton dynamics for $s_1 = 1.0$, $s_2 = 1.0$, $s = 1.0$

$$f(t, x, k, \omega) = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(kx - \omega t)$$

$$u(0, x) = f(0, x + 40, 0.2, 4 \cdot 0.2^3)$$



Soliton dynamics for $s_1 = 0.0$, $s_2 = 1.0$ $s = 1.0$

$$f(t, x, k, \omega) = \frac{\omega}{6k} + \frac{4}{3}k^2 - 2k^2 \tanh^2(kx - \omega t)$$

$$u(0, x) = f(0, x + 40, 0.2, 4 \cdot 0.2^3)$$

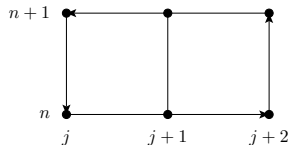
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System of KdV-Like Equations

$$\begin{cases} u_{1t} + 6u_1u_{1x} + u_{1xxx} - s_1u_1^2u_{1xx} - s_2u_{1xx} + u_1 - u_2 = 0, \\ u_{2t} + 6u_2u_{2x} + u_{2xxx} - s_1u_2^2u_{2xx} - s_2u_{2xx} + u_2 - u_1 = 0. \end{cases}$$

$$\begin{cases} \oint_{\partial\Omega} \left(-3u_1^2 + u_{1xx} - \frac{s_1}{3}u_1^3 - s_2u_{1x} \right) dt + u_1 dx + \iint_{\Omega} (u_1 - u_2) dt dx = 0, \\ \oint_{\partial\Omega} \left(-3u_2^2 + u_{2xx} - \frac{s_1}{3}u_2^3 - s_2u_{2x} \right) dt + u_2 dx + \iint_{\Omega} (u_2 - u_1) dt dx = 0. \end{cases}$$



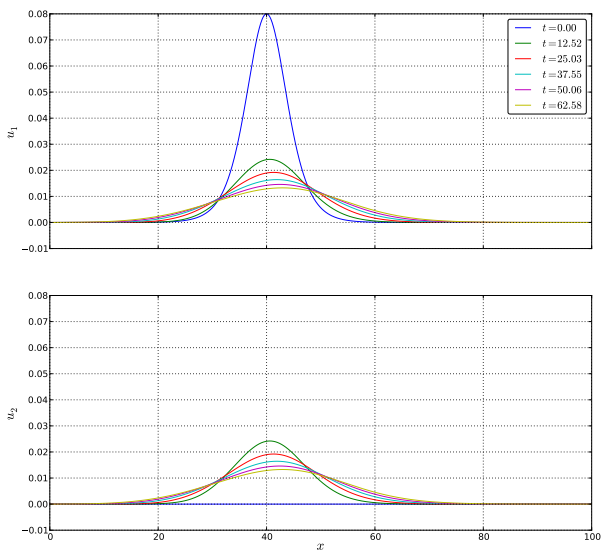
$$\int_{x_j}^{x_{j+1}} u_{1x} dx = u_1(t, x_{j+1}) - u_1(t, x_j), \quad \int_{x_j}^{x_{j+1}} u_{1xx} dx = u_{1x}(t, x_{j+1}) - u_{1x}(t, x_j),$$

$$\int_{x_j}^{x_{j+1}} u_{2x} dx = u_2(t, x_{j+1}) - u_2(t, x_j), \quad \int_{x_j}^{x_{j+1}} u_{2xx} dx = u_{2x}(t, x_{j+1}) - u_{2x}(t, x_j).$$

Discretization

$$\left\{ \begin{array}{l} -(1 + \theta_t - \theta_x^2 - \theta_t \theta_x^2) \circ (3u_1^2 + u_{1xx} - \frac{s_1}{3} u_1^3 - s_2 u_{1x}) \cdot \frac{\tau}{2} + \\ \quad + (\theta_x \theta_t - \theta_x) \circ u_1 \cdot 2h + (\theta_x \theta_t - \theta_x) \circ (u_1 - u_2) \cdot h\tau = 0, \\ -(1 + \theta_t - \theta_x^2 - \theta_t \theta_x^2) \circ (3u_2^2 + u_{2xx} - \frac{s_1}{3} u_2^3 - s_2 u_{2x}) \cdot \frac{\tau}{2} + \\ \quad + (\theta_x \theta_t - \theta_x) \circ u_2 \cdot 2h + (\theta_x \theta_t - \theta_x) \circ (u_2 - u_1) \cdot h\tau = 0, \\ (\theta_x + 1) \circ u_{1x} \cdot \frac{h}{2} = (\theta_x - 1) \circ u_1, \\ \theta_x \circ u_{1xx} \cdot 2h = (\theta_x^2 - 1) \circ u_{1x}, \\ (\theta_x + 1) \circ u_{2x} \cdot \frac{h}{2} = (\theta_x - 1) \circ u_2, \\ \theta_x \circ u_{2xx} \cdot 2h = (\theta_x^2 - 1) \circ u_{2x}. \end{array} \right.$$

$$\text{FDA} = \left\{ \begin{aligned}
 & \frac{u_1^{n+1} - u_1^n}{3} + \frac{(u_{1j+1}^{2n+1} - u_{1j-1}^{2n+1}) + (u_{1j+1}^{2n} - u_{1j-1}^{2n})}{4h} + \\
 & + \frac{(u_{1j+2}^{n+1} - 2u_{1j+1}^{n+1} + 2u_{1j-1}^{n+1} - u_{1j-2}^{n+1}) + (u_{1j+2}^n - 2u_{1j+1}^n + 2u_{1j-1}^n - u_{1j-2}^n)}{4h^3} - \\
 & - \frac{s_1 (u_{1j+1}^{3n+1} - u_{1j-1}^{3n+1}) + (u_{1j+1}^{3n} - u_{1j-1}^{3n})}{3} - \\
 & - s_2 \frac{(u_{1j+1}^{n+1} - 2u_{1j}^{n+1} + u_{1j-1}^{n+1}) + (u_{1j+1}^n - 2u_{1j}^n + u_{1j-1}^n)}{2h^2} + \\
 & + (u_1^{n+1} + u_1^n) \frac{1}{2} - (u_2^{n+1} + u_2^n) \frac{1}{2} = 0, \\
 & \frac{u_2^{n+1} - u_2^n}{3} + \frac{(u_{2j+1}^{2n+1} - u_{2j-1}^{2n+1}) + (u_{2j+1}^{2n} - u_{2j-1}^{2n})}{4h} + \\
 & + \frac{(u_{2j+2}^{n+1} - 2u_{2j+1}^{n+1} + 2u_{2j-1}^{n+1} - u_{2j-2}^{n+1}) + (u_{2j+2}^n - 2u_{2j+1}^n + 2u_{2j-1}^n - u_{2j-2}^n)}{4h^3} - \\
 & - \frac{s_1 (u_{2j+1}^{3n+1} - u_{2j-1}^{3n+1}) + (u_{2j+1}^{3n} - u_{2j-1}^{3n})}{3} - \\
 & - s_2 \frac{(u_{2j+1}^{n+1} - 2u_{2j}^{n+1} + u_{2j-1}^{n+1}) + (u_{2j+1}^n - 2u_{2j}^n + u_{2j-1}^n)}{2h^2} + \\
 & + (u_2^{n+1} + u_2^n) \frac{1}{2} - (u_1^{n+1} + u_1^n) \frac{1}{2} = 0.
 \end{aligned} \right.$$



Soliton dynamics for $s_1 = 0.0$, $s_2 = 1.0$

$$f(t, x, k) = 2k^2(1 - \tanh(kx - 4k^3t)^2)$$

$$u_1(0, x) = f(0, x + 40, 0.2), \quad u_2(0, x) = 0$$

Conclusions

- We applied our computer algebra assisted approach [1][Gerdt, Blinkov, Mozzhilkin'2006] based on finite volume method and difference elimination to obtain FDA to the KdV equation itself and to some KdV-like PDEs that model dynamics of elastic cylinder shells containing viscous incompressible liquid.
- The structure of a FDA depends on the numerical methods used to approximate integrals.
- All obtained FDAs are s-consistent.
- By numerical experiments we shown that the obtained FDAs nicely describe the dynamics of solitons and other exact solutions.
- Gröbner bases techniques can be used for difference elimination and for verification of s-consistency [3][Gerdt'2012].
- Difference Gröbner bases algorithms are described in [2][Gerdt, Robertz'2012], [3][Gerdt'2012] and [4][Scala'2011].

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