

A higher order Multiscale Finite Element Method for num. homy.

1

Let $\Omega \subset \mathbb{R}^n$ open, bounded, $f \in H^{-1}(\Omega)$.

$$(1) \quad \begin{cases} -\nabla \cdot (a^\varepsilon(x) \nabla u^\varepsilon) = f & \text{in } \Omega \\ u^\varepsilon = 0 & \text{on } \partial\Omega \end{cases}$$

$a^\varepsilon(x) = a(\frac{x}{\varepsilon})$, $a \in L^\infty(Y)$ with $Y = (0,1)^n$, a positive, Y -periodic

We know: For ε small, u^ε is approximated by

$$u^\varepsilon(x) \approx u^*(x) + \varepsilon \sum_{i=1}^n \chi_i(\frac{x}{\varepsilon}) \frac{\partial u^*}{\partial x_i}(x)$$

$$\text{and } \nabla u^\varepsilon(x) \approx \nabla u^*(x) + \sum_{i=1}^n (\nabla_y \chi_i)(\frac{x}{\varepsilon}) \frac{\partial u^*}{\partial x_i}(x)$$

where χ_i is the solution of the cell problem

$$\begin{cases} -\nabla_y \cdot (a(y)(e_i + \nabla_y \chi_i)) = 0 & \text{in } Y \\ y \rightarrow \chi_i(y) \text{ } Y\text{-periodic} \end{cases}$$

and u^* solution of the homogenized problem

$$\begin{cases} -\nabla \cdot (A^* \nabla u^*) = f & \text{in } \Omega \\ u^* = 0 & \text{on } \partial\Omega \end{cases}$$

$$\text{where } A^* e_i = \int_Y a(y)(e_i + \nabla_y \chi_i) dy$$

Idea: Approximation of u^ε of the form (1st. order Taylor expansion)

$$u^\varepsilon(x) \approx u^*(x + \varepsilon \chi(\frac{x}{\varepsilon})), \quad \chi = (\chi_1, \dots, \chi_n)^T$$

If $\{\phi_i^h\}_i$ classical FE-basis on coarse mesh of mesh size $h > \varepsilon$,
introduce oscillating basis (MsFEM) def. by (composition rule)

$$\phi_i^{\varepsilon,h}(x) = \phi_i^h(x + \varepsilon \chi(\frac{x}{\varepsilon}))$$

• change of variables: $x \rightarrow \tilde{w}^\varepsilon(x) := x + \varepsilon \chi(\frac{x}{\varepsilon})$. (Standard Galerkin Procedure to variational form of (1) will)

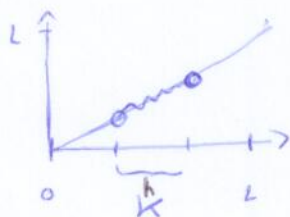
2) Definition of the MsFEM

let \mathcal{T}_h be a (coarse) conformal triangulation of Ω .

For $i = 1, \dots, n$ define $\hat{w}_i^{\varepsilon, h} \in H^1(\Omega)$ locally for each $K \in \mathcal{T}_h$ by

$$(2) \quad \begin{cases} -\nabla \cdot (a^\varepsilon \nabla \hat{w}_i^{\varepsilon, h}|_K) = -\nabla \cdot (A_K^* \nabla x_i) & \text{in } K \\ \hat{w}_i^{\varepsilon, h}|_K = x_i & \text{on } \partial K \end{cases}$$

and $\hat{w}^{\varepsilon, h} = (\hat{w}_1^{\varepsilon, h}, \dots, \hat{w}_n^{\varepsilon, h})^T \in H^1(\Omega, \mathbb{R}^n)$, A_K^* local average of a^ε



Remark: In the following, $\hat{w}_i^{\varepsilon, h}$ denotes the exact solution of (2),
while $w_i^{\varepsilon, h}$ denotes the approximated solution (\mathbb{P}_K -Lagrange-FE)
on a fine mesh $\mathcal{T}_h^{K'}$ of mesh size h' (oscillating test-f)

$$\underbrace{0 < h' < \varepsilon < h < 1}_{\mathbb{P}_K}$$

MsFEM - Basis:

let $V_h \subset H_0^1(\Omega)$ be the space of \mathbb{P}_K -Lagrange FEs on \mathcal{T}_h
with basis $\{\phi_i^h\}_{i=1}^{N_h}$. We define the MsFEM basis by

$$\phi_i^{\varepsilon, h}(x) = \phi_i^h \circ \omega^{\varepsilon, h}(x) \quad i = 1, \dots, N_h$$

Thus, $V_h^\varepsilon := \text{span} \{\phi_i^{\varepsilon, h}\}_{i=1}^{N_h}$ is conforming ($V_h^\varepsilon = V_h \circ \omega^{\varepsilon, h}$)

$$(3) \quad \begin{cases} \text{Find } u_h \circ \omega^{\varepsilon, h} \in V_h^\varepsilon : \\ a^\varepsilon(u_h \circ \omega^{\varepsilon, h}, v_h \circ \omega^{\varepsilon, h}) = \int_\Omega f v \circ \omega^{\varepsilon, h} dx \quad \forall v \circ \omega^{\varepsilon, h} \in V_h \end{cases}$$

Note that $V_h^\varepsilon = H_0^1 \circ \hat{w}^{\varepsilon, h} \subset H_0^1$ if $\hat{w}^{\varepsilon, h} \in W^{1, \infty}(\Omega, \mathbb{R}^n)$ (+dox Hilgrom, unique solution)

Example: \mathbb{P}_k -FES on \mathcal{T}_h ($k=1$)

p_i basis fct. of V_h

$$\Rightarrow \phi_i^{\varepsilon, h}(x) = p_i \circ \omega^{\varepsilon, h}(x)$$

$$\stackrel{\text{linear}}{=} p_i(x) + \sum_{j=1}^n (\omega_j^{\varepsilon, h}(x) - x_j) \frac{\partial p_i}{\partial x_j}(x)$$

$$\Rightarrow \nabla \cdot (a^\varepsilon \nabla \phi_i^{\varepsilon, h}) = \sum_{j=1}^n \nabla \cdot (a^\varepsilon \nabla \omega_j^{\varepsilon, h}) \frac{\partial p_i}{\partial x_j} \text{ in } K \in \mathcal{T}_h$$

choose $A_K^* = \text{const in } (2)$,

then we have for $K \in \mathcal{T}_h$

$$\begin{cases} \nabla \cdot (a^\varepsilon \nabla \phi_i^{\varepsilon, h}) = 0 & \text{in } K \\ \phi_i^{\varepsilon, h} = p_i & \text{on } \partial K \end{cases}$$

\leadsto coincides with the "standard" MsFEM (Hou, Wu)

Remark: For $k \geq 1$, let $n_j, j=1, \dots, N_h$ be the nodes associated to the \mathbb{P}_k -FEM on \mathcal{T}_h , $\{\phi_i^h\}_{i=1}^{N_h}$ Lagrange-basis of $V_h = \mathbb{P}_k$. It holds

$$\phi_i^h(n_j) = \delta_{ij}$$

$$\text{and } \phi_i^{\varepsilon, h}(n_j) = \delta_{ij} \quad \text{since } \omega^{\varepsilon, h}(x) \equiv x \text{ on } \partial K, K \in \mathcal{T}_h$$

Definition: For $v \in V_h = \mathbb{P}_k$ define the interpolation operators

$$\Pi_h v(x) := \sum_{e=1}^{N_h} v(n_e) \phi_e^h(x)$$

$$\Pi_h^\varepsilon v(x) := \sum_{e=1}^{N_h} v(n_e) \phi_e^{\varepsilon, h}(x) \quad (= (\Pi_h^\varepsilon v) \circ \omega^{\varepsilon, h}(x))$$

3) Convergence

4

Lemma 1:

Let $\omega \subset \mathbb{R}^n$ be open, bounded (smooth), $f \in L^2(\omega)$, $g \in H^1(\omega)$.

Let v^ε be the solution of (1) in ω , satisfying $v^\varepsilon = g$ on $\partial\omega$.

Let v^* be the homogenized limit and χ_i^ε the sol. of the cell problem.

Then there exists $C > 0$, indep. of $\varepsilon, \omega, f, g$:

$$\|v^\varepsilon - v^* - \varepsilon \sum_{i=1}^n \chi_i\left(\frac{x}{\varepsilon}\right) \frac{\partial v^*}{\partial x_i}(x)\|_{H^1(\omega)} \leq C \sqrt{\varepsilon} \sqrt{|\partial\omega|} \|v^*\|_{W^{2,\infty}(\omega)}$$

Lemma 2:

Let $\tilde{\omega}^\varepsilon(x) = x + \varepsilon \chi\left(\frac{x}{\varepsilon}\right) \in W^{1,\infty}(\Omega)$, $\chi_i \in W^{1,\infty}(\mathbb{R}^n)$, $u^* \in W^{2,\infty}(\Omega)$.

Then there exists $C \neq C(\varepsilon)$:

$$\|u^\varepsilon - u^* \circ \tilde{\omega}^\varepsilon\|_{H_0^1(\Omega)} \leq C \sqrt{\varepsilon}$$

Sketch of the proof:

$$\|\nabla u^\varepsilon - \nabla(u^* \circ \tilde{\omega}^\varepsilon)\|_{L^2(\Omega)^n}$$

$$\leq \|\nabla u^\varepsilon - (\nabla \tilde{\omega}^\varepsilon) \nabla u^*\|_{L^2(\Omega)^n}$$

$$\leq C \sqrt{\varepsilon} \text{ (Lemma 1)}$$

$$+ \|(\nabla \tilde{\omega}^\varepsilon) (\nabla u^* - (\nabla u^*) \circ \tilde{\omega}^\varepsilon)\|_{L^2(\Omega)^n}$$

$$\leq C \varepsilon \text{ (Taylor exp.)}$$

Theorem: Let u^ε be the exact solution of (1) and

$u_h^\varepsilon = u_h \circ \omega^{\varepsilon,h}$ be the numerical solution of (3).

Assume $u^* \in W^{k+1,\infty}(\Omega)$, $\chi_i \in W^{1,\infty}(\Omega)$, $\hat{\omega}^{\varepsilon,h} \in W^{1,\infty}(\Omega)$

Then there is a constant C indep. of ε and h such that

$$\|u^\varepsilon - u_h^\varepsilon\|_{H_0^1(\Omega)} \leq C \left(h^k + \sqrt{\frac{\varepsilon}{h}} + \left(\frac{h'}{\varepsilon}\right)^{k'} \right)$$

$k=1: (\varepsilon + h + \sqrt{\frac{\varepsilon}{h}} + \dots)$

Proof: From Céa's lemma, $\exists C$ indep. of ε, h s.t.

$$\|u^\varepsilon - u_h^\varepsilon\|_{H_0^1(\Omega)} \leq C \inf_{v_h^\varepsilon \in V_h^\varepsilon} \|u^\varepsilon - v_h^\varepsilon\|_{H_0^1(\Omega)}$$

Let $v_h^\varepsilon = \pi_h^\varepsilon u^*$, we have

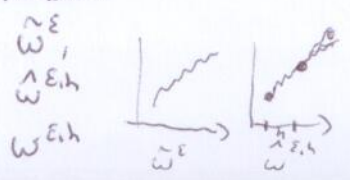
$$\|u^\varepsilon - u_h^\varepsilon\|_{H_0^1(\Omega)} \leq \|u^\varepsilon - \pi_h^\varepsilon u^*\|_{H_0^1(\Omega)}$$

Define $\tilde{\omega}^\varepsilon(x) := x + \varepsilon \chi(\frac{x}{\varepsilon})$, (rescaled sol. of the cell problem)

We get

$$\begin{aligned} & \|\nabla u^\varepsilon - \nabla(\pi_h^\varepsilon u^*)\|_{L^2(\Omega)^n} \\ & \leq C \|\nabla u^\varepsilon - \nabla(u^* \circ \tilde{\omega}^\varepsilon)\|_{L^2(\Omega)^n} \quad (\text{corrector result in periodic homog.}) \\ & \quad + \|\nabla(u^* - \pi_h u^*) \circ \tilde{\omega}^\varepsilon\|_{L^2(\Omega)^n} \quad (\text{interpolation result for the coarse mesh}) \\ & \quad + \|\nabla(\pi_h u^* \circ (\tilde{\omega}^\varepsilon - \hat{\omega}^{\varepsilon,h}))\|_{L^2(\Omega)^n} \quad (\text{homog. result of the local oscillating test}) \\ & \quad + \|\nabla(\pi_h u^* \circ (\hat{\omega}^{\varepsilon,h} - \omega^{\varepsilon,h}))\|_{L^2(\Omega)^n} \quad (\text{error estimate for the } P_k\text{-FEM used to compute the local oscillating test}) \end{aligned}$$

Recapitulate:



$$(i) \quad \|\nabla u^\varepsilon - \nabla u^* \circ \tilde{\omega}^\varepsilon\|_{L^2(\Omega)^n} \leq \underline{C\sqrt{\varepsilon}} \quad \text{by lemma 2}$$

$$(ii) \quad \|(\nabla \tilde{\omega}^\varepsilon)(\nabla(u^* - \Pi_h u^*) \circ \tilde{\omega}^\varepsilon)\|_{L^2(\Omega)^n}$$

$$\leq \|\text{Id} + \nabla_Y \chi\|_{L^\infty(Y)} \|\nabla(u^* - \Pi_h u^*) \circ \tilde{\omega}^\varepsilon\|_{L^2(\Omega)^n}$$

We have

$$\begin{aligned} \nabla(u^* - \Pi_h u^*) \circ \tilde{\omega}^\varepsilon &\stackrel{\text{Taylor Exp.}}{=} \nabla(u^* - \Pi_h u^*)(x) \\ &\quad + \varepsilon \int_0^1 \nabla \nabla(u^* - \Pi_h u^*)(x + \varepsilon t \chi) \cdot \chi dt \end{aligned}$$

$$\begin{aligned} \Rightarrow (ii) &\leq \|\text{Id} + \nabla_Y \chi\|_{L^\infty(Y)} \left(\|u^* - \Pi_h u^*\|_{H^1(\Omega)} + \varepsilon \|u^* - \Pi_h u^*\|_{W^{2,\infty}(\Omega)} \|\chi\|_{L^2} \right) \\ &\stackrel{\text{interp. result for degree } \leq \text{Interp.}}{\leq} \|\text{Id} + \nabla_Y \chi\|_{L^\infty(Y)} \|u^*\|_{W^{k+1,\infty}(\Omega)} \left(h^k + \varepsilon h^{k-1} \|\chi\|_{L^2(Y)} \right) \end{aligned}$$

$$(iii) \quad \|(\nabla \tilde{\omega}^\varepsilon - \nabla \hat{\omega}^{\varepsilon,h})(\nabla(\Pi_h u^*) \circ (\tilde{\omega}^\varepsilon - \hat{\omega}^{\varepsilon,h}))\|_{L^2(\Omega)^n}$$

$$\leq \|u^*\|_{W^{1,\infty}(\Omega)}^* \|\nabla(\tilde{\omega}^\varepsilon - \hat{\omega}^{\varepsilon,h})\|_{L^2(\Omega)^n}$$

It holds

$$\begin{aligned} \|\nabla(\tilde{\omega}^\varepsilon - \hat{\omega}^{\varepsilon,h})\|_{L^2(\Omega)^n}^2 &= \sum_{k \in \mathcal{J}_h} \|\nabla(\tilde{\omega}^\varepsilon - \hat{\omega}^{\varepsilon,h})\|_{L^2(\Omega)}^2 \\ &\stackrel{\text{lemma 1}}{\leq} C \sum_{k \in \mathcal{J}_h} \varepsilon h^{n-1} \stackrel{\#k=h^{-n}}{=} C \varepsilon \frac{h^{n-n}}{h^n} = \underline{C \frac{\varepsilon}{h}} \end{aligned}$$

$$(iv) \quad \|\nabla(\hat{\omega}^{\varepsilon,h} - \omega^{\varepsilon,h})(\nabla \Pi_h u^*) \circ (\hat{\omega}^{\varepsilon,h} - \omega^{\varepsilon,h})\|_{L^2(\Omega)^n}$$

$$\leq \|u^*\|_{W^{1,\infty}(\Omega)} \|\nabla \hat{\omega}^{\varepsilon,h} - \omega^{\varepsilon,h}\|_{L^2(\Omega)^n}$$

$$\begin{aligned} \text{We have } \|\nabla(\hat{\omega}^{\varepsilon,h} - \omega^{\varepsilon,h})\|_{L^2(\Omega)^n}^2 &= \sum_{k \in \mathcal{J}_h} \|\nabla(\hat{\omega}^{\varepsilon,h} - \omega^{\varepsilon,h})\|_{L^2(k)}^2 \\ &\stackrel{\mathbb{P}_h\text{-interp.}}{\leq} C \sum_{k \in \mathcal{J}_h} (h')^{2k'} \underbrace{|\hat{\omega}^{\varepsilon,h}|_{H^{k'+1}(k)}^2}_{\leq (C \varepsilon^{-k'} \sqrt{|k|})^2} \quad \leadsto (iv) \leq C \varepsilon^{-k'} h^{k'} \frac{h^{\frac{n}{2}}}{h^{\frac{n}{2}}} \\ &\leq \end{aligned}$$

$$(i)-(iv) \Rightarrow \|\nabla u^\varepsilon - \nabla(\Pi_h^\varepsilon u^*)\| \leq C \left(\sqrt{\varepsilon} + (h^k + \varepsilon h^{k-1}) + \sqrt{\frac{\varepsilon}{h}} + \left(\frac{h'}{\varepsilon}\right)^{k'} \right) \quad \square$$